Measurement and Control of the Longitudinal Phase Space at High-Gain Free-Electron Lasers.

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Deutsches Elektronen-Synchrotron (DESY)

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- Introduction and Motivation
- Principle of Longitudinal Phase Space Diagnostics
- 3 Control and Manipulation of the Longitudinal Phase Space
- Special Applications of Longitudinal Phase Space Diagnostics
- 5 Summary and Conclusions

Motivation for Longitudinal Phase Space Diagnostics

Requirements of high-gain free-electron lasers including various seeded schemes

- ▶ Good electron beam quality in terms of energy spread, emittance, and peak current
- ➤ Time-resolved (i.e. longitudinal) information and control of these parameters
- ▶ Control of beam instabilities and corresponding diagnostics problems, e.g. COTR
- Measurement and tunability of electron and photon pulse lengths

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Measurement and control of the longitudinal phase space (t, δ)

- Longitudinal position t=-z/c and relative energy or momentum deviation $\delta=\frac{\Delta p}{p_0}$
- 6-d beam transport matrix provide a powerful formalism to discuss the underlying physics

$$\begin{pmatrix} x_f \\ x_f' \\ y_f \\ y_f' \\ t_f \\ \delta_f \end{pmatrix} = \begin{pmatrix} R_{11} & R_{12} & R_{13} & R_{14} & R_{15} & R_{16} \\ R_{21} & R_{22} & R_{23} & R_{24} & R_{25} & R_{26} \\ R_{31} & R_{32} & R_{33} & R_{34} & R_{35} & R_{36} \\ R_{41} & R_{42} & R_{43} & R_{44} & R_{45} & R_{46} \\ R_{51} & R_{52} & R_{53} & R_{54} & R_{55} & R_{56} \\ R_{61} & R_{62} & R_{63} & R_{64} & R_{65} & R_{66} \end{pmatrix} \cdot \begin{pmatrix} x_i \\ x_i' \\ y_i \\ y_i' \\ t_i \\ \delta_i \end{pmatrix}$$

Motivation for Longitudinal Phase Space Diagnostics

Requirements of high-gain free-electron lasers including various seeded schemes

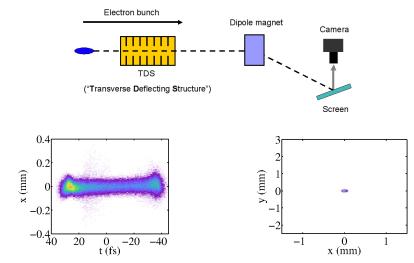
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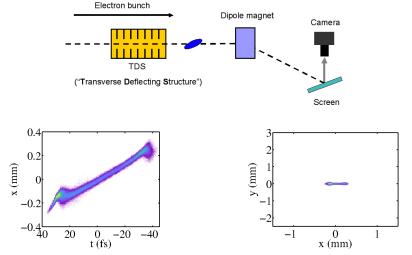
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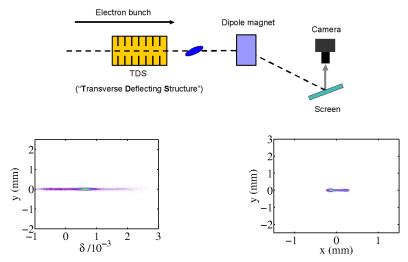
- ★ Transverse Deflecting RF Structure in combination with an energy (dipole) spectrometer
- Single-shot capability and high resolution in both energy and time



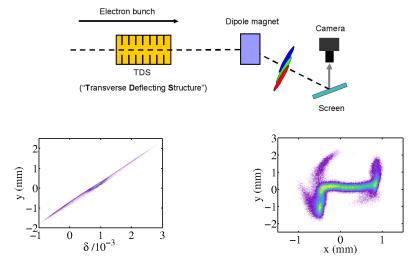
Start with initial distributions in (t, x) and (x, y)



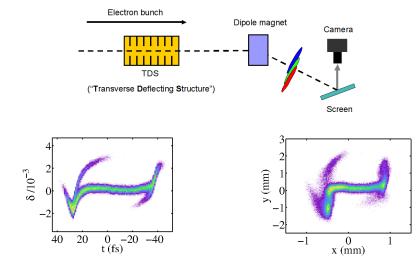
- TDS imposes a time-dependent transversal kick $\Delta x'(t) \sim \sin(t) \approx t$ (at zero-crossing)
- Appropriate beam transport optics (R_{12}) maps $\Delta x'(t) \to \Delta x(t)$, i.e. $\Delta x(t) \sim t$
- Time information is translated to the horizontal position



▶ Start again with present distributions in (δ, y) and (x, y)



- ▶ Dipole magnet imposes a energy-dependent transversal kick $\Delta y'(\delta) \sim \delta$
- ▶ Appropriate beam transport optics (R_{34}) maps $\Delta y'(\delta) \to \Delta y(\delta)$, i.e. $\Delta y(\delta) \sim \delta$
- ★ Energy information is translated to the vertical position



- **\star** Transformation of the longitudinal phases space (t, δ) to (x, y)
- ★ Good agreement compared to the real longitudinal phase space
- ★ Simulation shows some discrepancy when looking into the details

Measureable beam size and time resolution using a TDS

$$au$$
 $x(s,t)=x_0(s)+S(s)\cdot t$ with shear function $S(s)=\frac{eV_0\omega}{E}\sqrt{\beta_1\beta_2(s)}\sin(\Delta\Psi_x(s))$ ($S=R_{15}$)

$$lacktriangledown$$
 $\sigma_{x}=\sqrt{\sigma_{x_0}^2+(S\cdot\sigma_t)^2}\Rightarrow$ r.m.s. definition of time resolution $\sigma_{R,t}=\sigma_{x_0}/S$

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Measureable beam size and energy resolution using an energy spectrometer

- $y(s, \delta) = y_0(s) + D(s) \cdot \delta$ with dispersion function D(s) ($D = R_{36}$)
- lacktriangledown $\sigma_{
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- \star Small intrinsic beam sizes at screen position and large S and D improve resolution

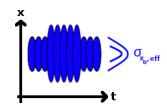
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- \star Small intrinsic beam sizes at screen position and large S and D improve resolution
- ★ Intrinsic beam size could vary along the bunch



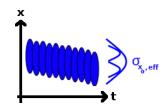
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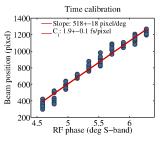
Measureable beam size and energy resolution using an energy spectrometer

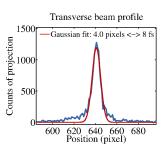
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- \star Small intrinsic beam sizes at screen position and large S and D improve resolution
- ★ Intrinsic beam size could vary along the bunch
- ★ Bunch could have a tilt
- ⇒ Definition describes the overall resolution
- The same is valid for the energy resolution



Calibration, Resolution, and Impact of Jitter: Time

Time calibration: Scanning of RF phase

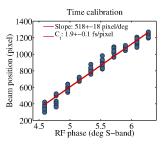


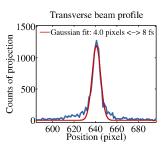


- ▶ Time calibration: Centroid offset versus RF phase ($\phi = \omega t \sim t$)
- ▶ Time resolution: Beam size without shearing by the TDS

Calibration, Resolution, and Impact of Jitter: Time

Time calibration: Scanning of RF phase





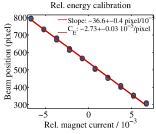
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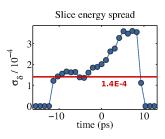
Impact of jitter sources on tranverse jitter

- Might be a problem for calibration which is a multi-shot procedure
- The only relevant jitter is arrival time σ_t and RF phase jitter σ_ϕ $\to \sigma_x = S \cdot \sigma_t$ and $\sigma_x = S \cdot \omega^{-1} \cdot \sigma_\phi$
- Basically this is under control for stable machines and can even be improved

Calibration, Resolution, and Impact of Jitter: Energy

Energy calibration: Scanning of magnet current (energy)

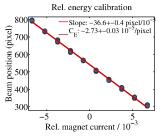


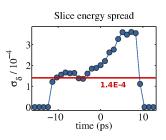


- ▶ Energy calibration: Centroid offset versus energy or simply magnet current ($\Delta I/I_0$)
- ▶ Energy resolution: Minimum slice energy spread of uncompressed bunches (reso. limited)

Calibration, Resolution, and Impact of Jitter: Energy

Energy calibration: Scanning of magnet current (energy)





- ▶ Energy calibration: Centroid offset versus energy or simply magnet current ($\Delta I/I_0$)
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Impact of jitter sources on tranverse jitter

- Might be a problem for calibration which is a multishot procedure
- The only relevant jitter is energy jitter σ_{δ} $\rightarrow \sigma_{v} = D \cdot \sigma_{\delta}$
- ★ Basically this is under control, but pay attention to hysteresis effects

TDS-induced energy spread and chirp: Theory and Experiment

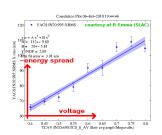
- Transverse deflecting structures induce
- Thick-lens matrix ⇒ induced energy chirp

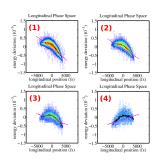
Transverse deflecting structures induce energy spread (Panofsky-Wenzel theorem) Thin-lens matrix
$$\rightarrow \sigma_{\delta} = K\sigma_{x} = \frac{eVk}{pc}\sigma_{x}$$
Thick-lens matrix \Rightarrow induced energy chirp

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- Thick-lens matrix ⇒ induced energy chirp

$$\begin{pmatrix} x_f \\ x'_f \\ t_f \\ \delta_f \end{pmatrix} = \begin{pmatrix} 1 & L & KL/2 & 0 \\ 0 & 1 & K & 0 \\ 0 & 0 & 1 & 0 \\ K & KL/2 & K^2L/6 & 1 \end{pmatrix} \cdot \begin{pmatrix} x_i \\ x'_i \\ t_i \\ \delta_i \end{pmatrix}$$

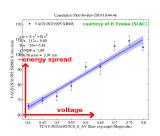


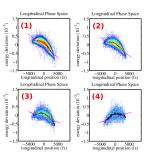


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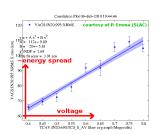
Coherent synchrotron radiation effects in the energy spectrometer

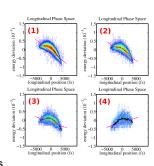
- ★ Basically negligible due to the fact that
 - time is transformed (e.g. in x) in front of the spectrometer
 - dispersion starts energy transformation (e.g. in y) before CSR is built up

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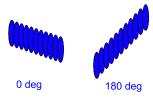




Wakefield effects

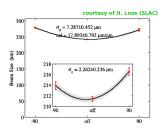
★ Basically negligible when having no large position offsets

Initial correlations in (x^\prime,t) may give different results when changing zero-crossing



•
$$\sigma_{x} = \sqrt{\sigma_{x_0}^2 + (C \pm S)^2 \cdot \sigma_t^2}$$

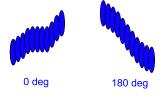
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★ If C is a constant: simple calculation using values at ±S (0 and 180 deg)

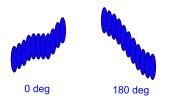
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- ★ If C varies along the bunch (i.e. C(t)): reconstruction from both projections is possible (idea and Ref. by H. Loos (SLAC))

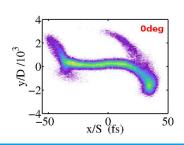
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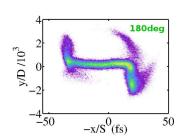


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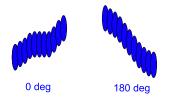
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Simulated measurements with both zero-crossings (0 and 180 deg)





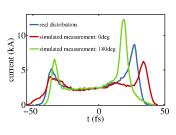
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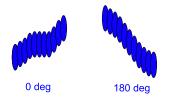
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· Strong effects in head and tail

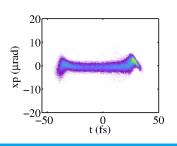
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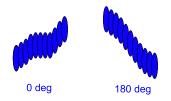
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Simulated measurements with both zero-crossings (0 and 180 deg)



- · Strong effects in head and tail
- ★ Linear scaling will not help

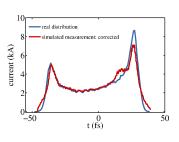
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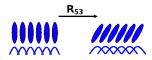
- · Strong effects in head and tail
- Linear scaling will not help
- ★ Reconstruction from two projections

Time coordinate along the energy spectrometer

▶ Longitudinal position after the spectrometer: $t_f = R_{53} \cdot y_i + R_{54} \cdot y_i' + t_i + R_{56} \cdot \delta_i$

Time coordinate along the energy spectrometer

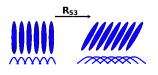
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- lacktriangledown $\sigma_{t_f}pprox R_{53}\cdot\sigma_{y_i}$ and $R_{53}pprox lpha$ i.e. the bending angle of the dipole
- ▶ $R_{53} \cdot \sigma_{y_i}$ smears out density modulations



Microbunches overlap and smear out

Time coordinate along the energy spectrometer

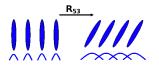
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- Microbunches overlap and smear out
- Larger beam size \rightarrow larger overlap

Time coordinate along the energy spectrometer

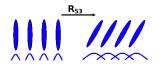
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- Microbunches overlap and smear out
- $\bullet \ \ \text{Larger beam size} \to \text{larger overlap}$
- Larger period → less overlap

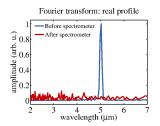
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- ▶ $R_{53} \cdot \sigma_{y_i}$ smears out density modulations



- Microbunches overlap and smear out
- Larger beam size → larger overlap
- Larger period → less overlap

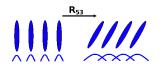
Particle tracking simulation with initial density modulation



- ▶ $5 \,\mu\mathrm{m}$ modulation in $(t, \delta)^{\mathrm{T}}$ before the spectrometer
- ▶ No modulation in $(t, \delta)^T$ after the spectrometer

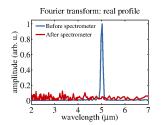
Time coordinate along the energy spectrometer

- ▶ Longitudinal position after the spectrometer: $t_f = R_{53} \cdot y_i + R_{54} \cdot y_i' + t_i + R_{56} \cdot \delta_i$
- $ightharpoonup \sigma_{t_f} pprox R_{53} \cdot \sigma_{y_i}$ and $R_{53} pprox \alpha$ i.e. the bending angle of the dipole
- ▶ $R_{53} \cdot \sigma_{y_i}$ smears out density modulations



- Microbunches overlap and smear out
- Larger beam size → larger overlap
- Larger period → less overlap

Particle tracking simulation with initial density modulation

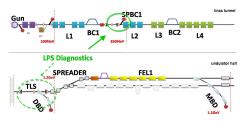


- ▶ 5 μm modulation in $(t, \delta)^{\text{T}}$ before the spectrometer
- ▶ No modulation in $(t, \delta)^T$ after the spectrometer
- \star Wavelengths $\lambda_c \ll 2\pi R_{51}\sigma_x$ will be suppressed
- ★ Proposal: Strong COTR mitigation in an energy spectrometer ⇒ emittance measurements

Longitudinal Phase Space Diagnostics at FERMI@Elettra

Courtesy of P. Craievich. For details: WEPA03.

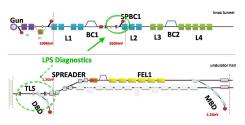
Free Electron Laser for Multidisciplinary Investigations (FERMI)



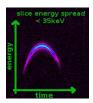
Longitudinal Phase Space Diagnostics at FERMI@Elettra

Courtesy of P. Craievich. For details: WEPA03.

Free Electron Laser for Multidisciplinary Investigations (FERMI)



Longitudinal phase space diagnostics at FERMI@Elettra

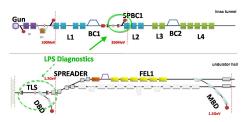


LPS data in SPBC1 at low energy

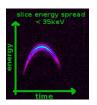
Longitudinal Phase Space Diagnostics at FERMI@Elettra

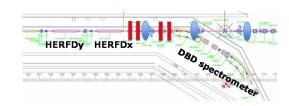
Courtesy of P. Craievich. For details: WEPA03.

Free Electron Laser for Multidisciplinary Investigations (FERMI)



Longitudinal phase space diagnostics at FERMI@Elettra





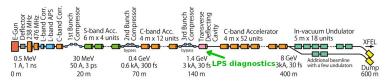
► LPS data in SPBC1 at low energy

High energy TDS (both planes) will be installed soon

Longitudinal Phase Space Diagnostics at SACLA/SPring-8

Courtesy of Y. Otake.

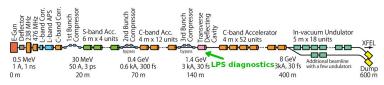
SPring-8 Angstrom Compact Free-Electron Laser (SACLA)



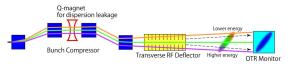
Longitudinal Phase Space Diagnostics at SACLA/SPring-8

Courtesy of Y. Otake.

SPring-8 Angstrom Compact Free-Electron Laser (SACLA)



Longitudinal phase space diagnostics at SACLA

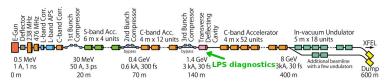


- Quadrupole kicks in a dispersive section
- Residual dispersion after the bunch compressor

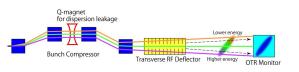
Longitudinal Phase Space Diagnostics at SACLA/SPring-8

Courtesy of Y. Otake.

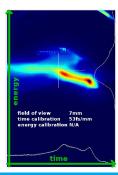
SPring-8 Angstrom Compact Free-Electron Laser (SACLA)



Longitudinal phase space diagnostics at SACLA

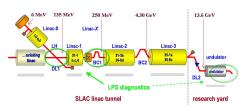


- Quadrupole kicks in a dispersive section
- Residual dispersion after the bunch compressor
- Still in commissioning phase
- Preliminary longitudinal phase space measurement



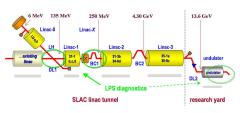
Courtesy of Y. Ding, P. Emma, and H. Loos.

The Linac Coherent Light Source (LCLS)

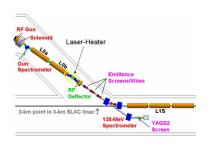


Courtesy of Y. Ding, P. Emma, and H. Loos.

The Linac Coherent Light Source (LCLS)



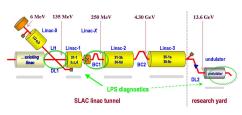
Longitudinal phase space diagnostics at LCLS



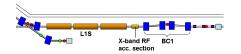
- Longitudinal phase space diagnostics
- Longitudinal phase space manipulation
 - Laser heater

Courtesy of Y. Ding, P. Emma, and H. Loos.

The Linac Coherent Light Source (LCLS)



Longitudinal phase space diagnostics at LCLS

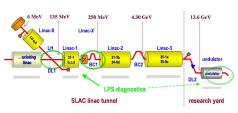


- Longitudinal phase space diagnostics
 - Longitudinal phase space manipulation

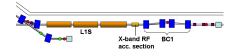
 Laser heater
- Longitudinal phase space linearization
 - X-band RF linearizer

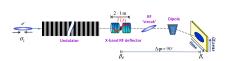
Courtesy of Y. Ding, P. Emma, and H. Loos.

The Linac Coherent Light Source (LCLS)



Longitudinal phase space diagnostics at LCLS

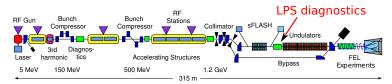




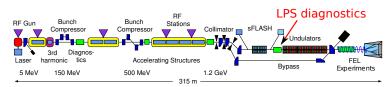
- Longitudinal phase space diagnostics
- Longitudinal phase space manipulation

 Laser heater
- Longitudinal phase space linearization
 X-band RF linearizer
- A balla IXI lillice
- In preparation
 X-band TDS after the undulators (project started this July)

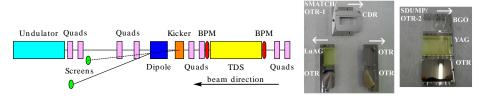
The Free-Electron Laser in Hamburg (FLASH)



The Free-Electron Laser in Hamburg (FLASH)

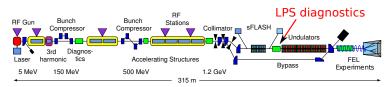


Longitudinal phase space diagnostics at FLASH

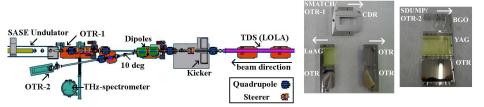


Longitudinal phase space diagnostics in front of the undulators

The Free-Electron Laser in Hamburg (FLASH)

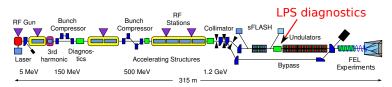


Longitudinal phase space diagnostics at FLASH

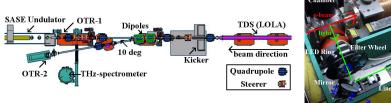


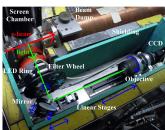
Longitudinal phase space diagnostics in front of the undulators

The Free-Electron Laser in Hamburg (FLASH)



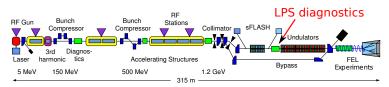
Longitudinal phase space diagnostics at FLASH



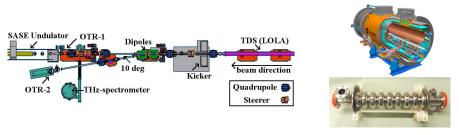


- Longitudinal phase space diagnostics in front of the undulators
- Indispensable for beam dynamics studies in general

The Free-Electron Laser in Hamburg (FLASH)

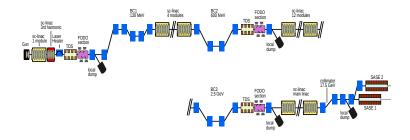


Longitudinal phase space diagnostics at FLASH

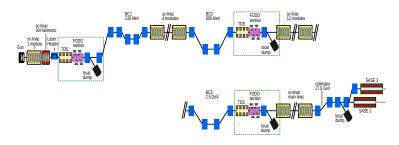


- Longitudinal phase space diagnostics in front of the undulators
- ▶ Indispensable for beam dynamics studies in general
- ▶ Longitudinal phase space linearizations with third-harmonic RF linearizer (3.9 GHz)

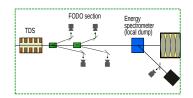
European X-Ray Free Electron Laser (E-XFEL)



European X-Ray Free Electron Laser (E-XFEL)



Longitudinal phase space diagnostics at E-XFEL

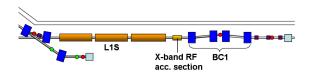


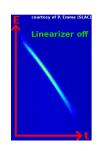
- Three setups for longitudinal phase space measurements (including slice emittance)
- Higher-harmonic RF linearizer and a laser heater
- Proposal: Apply longitudinal phase space diagnostics on individuals bunches of the train (septum magnet)

Longitudinal Phase Space Linearization at LCLS/SLAC

For details: TUOCAB02 by P. Emma et al. in the Proceedings of PAC'07

Linearization of the longitudinal phase space using an X-band RF linearizer





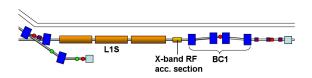
Measurements

Longitudinal phase space (off-crest):
 X-band linearizer switched off

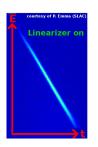
Longitudinal Phase Space Linearization at LCLS/SLAC

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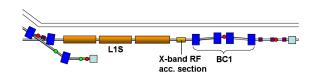
Measurements

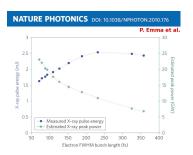
- Longitudinal phase space (off-crest):
 X-band linearizer switched off
- Longitudinal phase space (off-crest):
 X-band linearizer switched on

Longitudinal Phase Space Linearization at LCLS/SLAC

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Linearization of the longitudinal phase space using an X-band RF linearizer



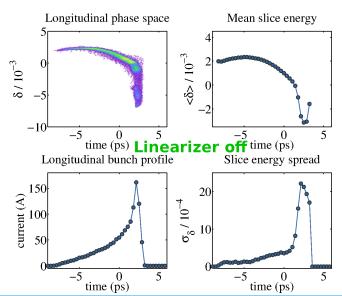


Measurements

- Longitudinal phase space (off-crest):
 X-band linearizer switched off
- Longitudinal phase space (off-crest):
 X-band linearizer switched on
- Control of the bunch lengths allows control of the FEL photon pulse durations (still a hot topic)

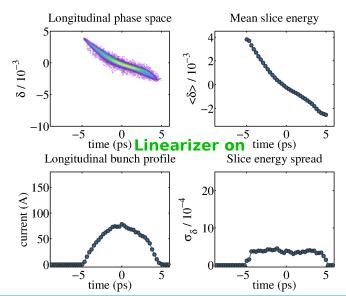
Longitudinal Phase Space Linearization at FLASH/DESY

Measurement of the longitudinal phase space: third-harmonic RF linearizer off



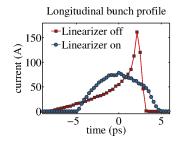
Longitudinal Phase Space Linearization at FLASH/DESY

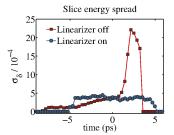
Measurement of the longitudinal phase space: third-harmonic RF linearizer on

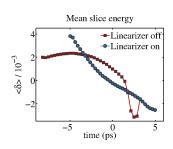


Longitudinal Phase Space Linearization at FLASH/DESY

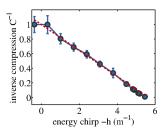
Comparison: third-harmonic RF linearizer on/off



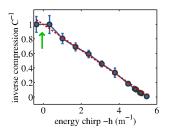


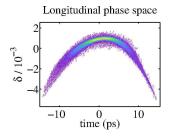


- ► Compression in a magnetic bunch compressor: $C^{-1} = (1 + h_1 R_{56}) + (h_2 R_{56} + 2h_1^2 T_{566})t_i$
- h_1,h_2 : first and second order energy chirp
- R₅₆, T₅₆₆: first and second oder longitudinal dispersion
- ▶ Eliminate time-dependency of C^{-1} by using a proper $h_2 \Rightarrow$ higher-harmonic RF system (dual-frequency)

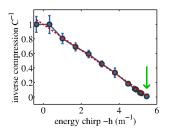


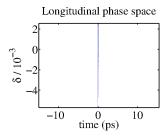
Measurements on linear bunch compression using only one bunch compressor



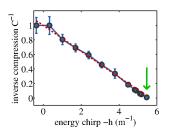


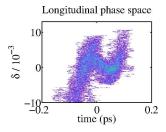
• Start with uncompressed bunches $(C^{-1}=1)$



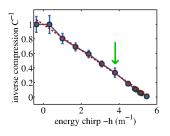


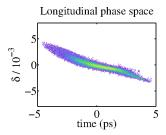
- Start with uncompressed bunches (C⁻¹ = 1)
- $\bullet~$ End with compressed bunches $({\it C}^{-1} \rightarrow 0)$



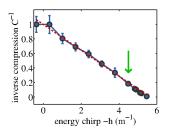


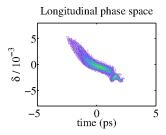
- Start with uncompressed bunches (C⁻¹ = 1)
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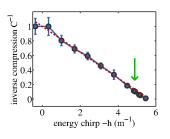


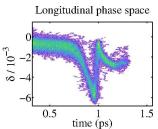
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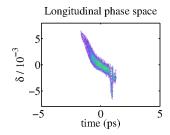




- Start with uncompressed bunches (C⁻¹ = 1)
- End with compressed bunches $(C^{-1} \rightarrow 0)$

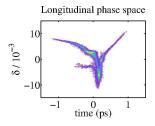






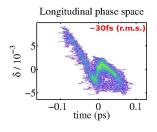
- Start with uncompressed bunches $(C^{-1} = 1)$
- End with compressed bunches $(C^{-1} \rightarrow 0)$
- ★ Strong local compression due to collective effects (not fully understood yet)

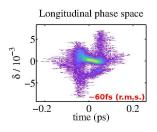
The old non-linear compression mode

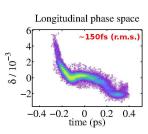


★ Non-linear compression: sharp leading spike (small charge fraction) with a long trailing tail

FEL operation with new linear compression mode

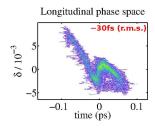


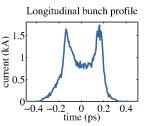


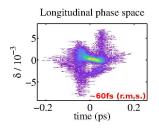


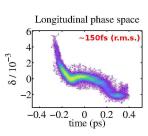
- ★ Non-linear compression: sharp leading spike (small charge fraction) with a long trailing tail
- ★ Linear compression: flexible bunch lengths and shapes (more regular but still complex)
- ★ Linear compression: more FEL pulse energies (at least 4×)

FEL operation with new linear compression mode



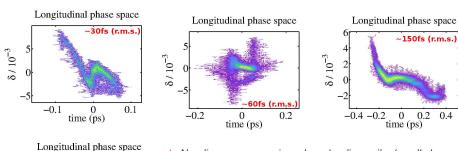


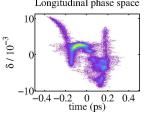




- ★ Non-linear compression: sharp leading spike (small charge fraction) with a long trailing tail
- ★ Linear compression: flexible bunch lengths and shapes (more regular but still complex)
- ★ Linear compression: more FEL pulse energies (at least 4×)
 - Appearance of double-horns like at LCLS

FEL operation with new linear compression mode





- ★ Non-linear compression: sharp leading spike (small charge fraction) with a long trailing tail
- ★ Linear compression: flexible bunch lengths and shapes (more regular but still complex)
- ★ Linear compression: more FEL pulse energies (at least 4×)
- Appearance of double-horns like at LCLS
- ★ Double-horns show some fragmentation

Tailoring the Longitudinal Phase Space for Wakefield Experiments

For details: P. Piot et al., Fermilab preprint PUB-11-339-APC (2011)

Beam-driven acceleration with drive and witness bunch

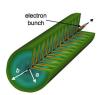


- Wakefield acceleration in dielectric structures using drive and witness bunches
- ▶ Transformer ratio $R = \frac{E_+}{E_-}$ is limited to ≤ 2 for symmetric current profiles

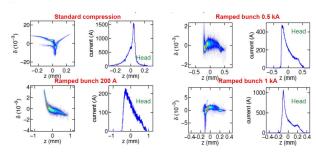
Tailoring the Longitudinal Phase Space for Wakefield Experiments

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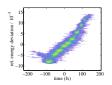
- Wakefield acceleration in dielectric structures using drive and witness bunches
- ▶ Transformer ratio $R = \frac{E_+}{E_-}$ is limited to ≤ 2 for symmetric current profiles
- ▶ Enhancement of the R by linearly ramped current profiles

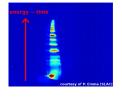


 \star Linearly ramped current profiles (from experiment) enables (from simulations) R > 6

Observations of Microbunching Instabilities in Time-domain

Linearly chirped bunches with intensity modulations at FLASH/DESY and LCLS/SLAC

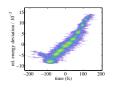


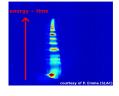


- Indication of microbunches
- Density modulations (FLASH)
- Energy modulations (LCLS)
 - ullet chirped bunches $o E \propto t$

Observations of Microbunching Instabilities in Time-domain

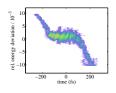
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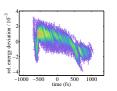




- Indication of microbunches
- Density modulations (FLASH)
- Energy modulations (LCLS)
 - chirped bunches \rightarrow $E \propto t$

Tilted microbunches and energy spread increase at FLASH/DESY

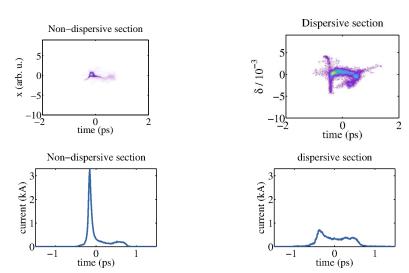




- Indication and observation of:
 - tilted microbunches
 - increased slice energy spread
- Common observation since operation with linear compression

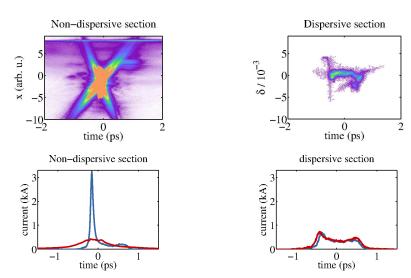
★Better control of slice energy spread is needed, especially for seeded FELs

Mitigation of Coherent Optical Transition Radiation: Experiment Compression instability: Charge 0.4 nC



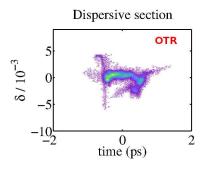
★ Strong discrepancy of current profiles between non-dispersive and dispersive section

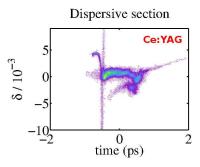
Mitigation of Coherent Optical Transition Radiation: Experiment Compression instability: Charge 0.5 nC



★ Local COTR emission spoils current profile measurement in non-dispersive section

Mitigation of Coherent Optical Transition Radiation: Experiment





- ★ COTR is most probably generated by an ultra-short local spike
- ★ No indication for COTR in the dispersive section
- Poster on mitigation of COTR in the non-dispersive section at FLASH (M. Yan et al. THPB16)

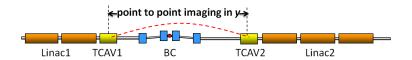
- Transverse deflecting structures induce energy spread (Panofsky-Wenzel theorem)
 ★ beam heating (cf. laser-heater)
- ▶ Energy spread is correlated, i.e. reversible ★ heat only where it's necessary

$$\begin{pmatrix} y_f \\ y_f' \\ t_f \\ \delta_f \end{pmatrix} = \begin{pmatrix} 1 & L & KL/2 & 0 \\ 0 & 1 & K & 0 \\ 0 & 0 & 1 & 0 \\ K & KL/2 & K^2L/6 & 1 \end{pmatrix} \cdot \begin{pmatrix} y_i \\ y_i' \\ t_i \\ \delta_i \end{pmatrix}$$

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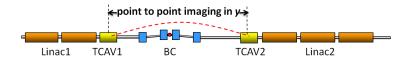
Reversible beam heating combining two transverse deflecting structures (cavities)



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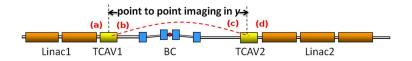
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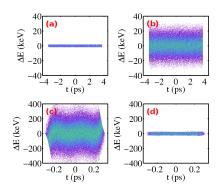


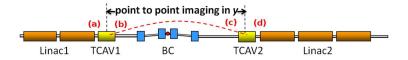
- ► Transport matrix: TCAV1 → TCAV2 (streaking ⊥ bending plane)
- ★ *K*₁*R*₅₆: effective energy spread for microbunching suppression
- energy spread and spatial chirp cancelation

$$\begin{pmatrix} a + \frac{bL_2}{2} + \frac{K_1K_2R_{56}L_2}{2} & \frac{L_2}{2a} & 0 & \frac{K_2L_2R_{56}}{2} \\ b + K_1K_2R_{56} & a^{-1} & 0 & K_2R_{56} \\ K_1R_{56} & 0 & 1 + hR_{56} & R_{56} \\ 0 & 0 & 0 & \frac{1}{1 + hR_{56}} \end{pmatrix}$$

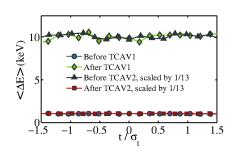


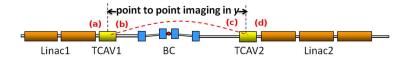
- Longitudinal phase space:
 - start with 1 keV slice energy spread
 - compression $C \approx 13$



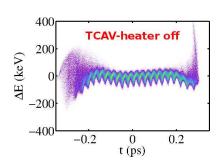


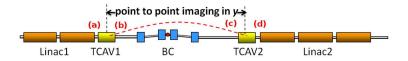
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- ★ Perfect cancelation of additional energy spread induced by TCAV1
 - CSR: small differences in the tails



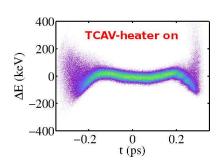


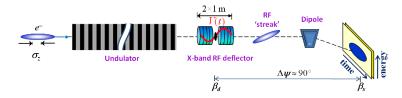
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- ★ TCAV-heater switched off

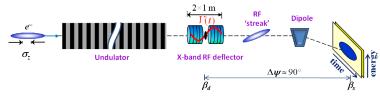




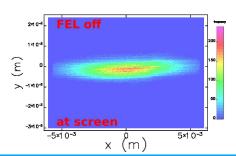
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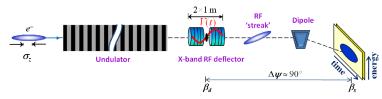




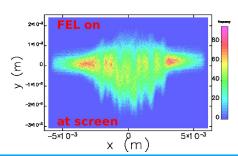


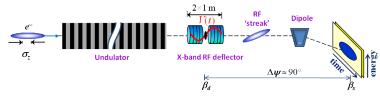
FEL: Genesis -> Tracking: elegant



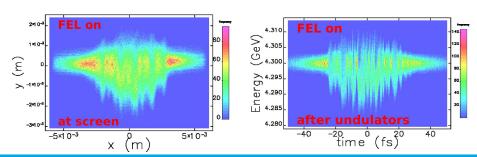


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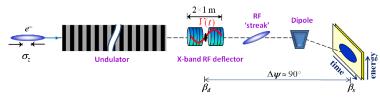




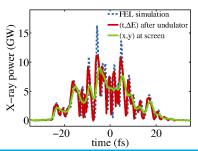
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Longitudinal phase space diagnostics right after the undulators

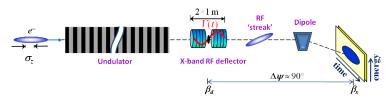


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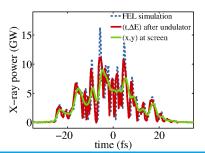


- Time-dependent energy loss and spread due to FEL process
- Correlation with temporal FEL photon pulse profile (replica of FEL photon pulse)

Longitudinal phase space diagnostics right after the undulators



FEL: Genesis -> Tracking: elegant



- Time-dependent energy loss and spread due to FEL process
- Correlation with temporal FEL photon pulse profile (replica of FEL photon pulse)
- ★ All features of high-resolution longitudinal phase space diagnostics (electrons beams):
 - independent of FEL wavelength
 - high-dynamic range
 - single-shot temporal profiles

Longitudinal phase space diagnostics based on TDS and energy spectrometer

- ★ Provide useful information on electron beams
- ★ Both with high resolution, high dynamic range, and single-shot capability

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Control of the longitudinal phase space by dual-frequency linear accelerators

- ★ Higher harmonic RF systems for longitudinal phase space linearizations

 linear compression
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Special applications of longitudinal phase space diagnostics

- ★ May provide useful information on X-ray pulses
- ★ Reversible Electron Beam Heater for Suppression of Microbunching Instabilities

Acknowledgments

Thanks for lots of helpful information and discussions to ...

- P. Evtushenko at Jlab
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